

Fig. 8. Vue perspective de la structure suivant y.

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## Crystal Structure of a $2M_2$ Lepidolite

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The crystal structure of a natural sample of a  $2M_2$  lepidolite from the Island of Elba, Tuscany (Italy), has been determined and refined by least-squares methods using three-dimensional data collected by the Weissenberg method with Cu  $K\alpha$  radiation. This specimen has a composition close to  $Pl_{50}Tl_{50}$  (where Pl = polyolithionite, Tl = trilithionite) and cell dimensions:  $a = 9.04 \pm 0.02$ ,  $b = 5.22 \pm 0.02$ ,  $c = 20.210 \pm 0.001$  Å,  $\beta = 99^\circ 35' \pm 20'$ . The main features of the structure are (a) a remarkable octahedral ordering, with one of the octahedral sites filled almost exclusively by Li atoms; (b) the tetrahedral sheets, made up by tetrahedra distorted into a form near to an elongated trigonal pyramid, deviate little from hexagonality ( $\alpha = 6^\circ 27'$ ).

### Introduction

The polymorphism of micas has long been, and still is, as subject of great interest to mineralogists ever since

Hendricks & Jefferson (1939) showed that these minerals crystallize with one of several layered structures, all of which are based on the same substructure. Many studies have been devoted to the problem of enumerat-

ing all possible stacking sequences of mica polymorphs, to the experimental methods of identification of such polymorphs, and to the survey of their frequencies of occurrence.

The present trend of research in the field of mica polymorphism is toward a more detailed knowledge of the structure of several polymorphs as an indispensable basis for understanding the factors which control their crystallization.

Following this trend we decided to undertake the determination of the structure of a  $2M_2$  lepidolite, whose features should be very important for the evaluation of the structural controls over mica polymorphism. The crystal structure of a  $2M_2$  lepidolite was published by Takeda, Haga & Sadanaga (1971) when the present determination was completed and the refinement was in process. As the chemical composition of this  $2M_2$  lepidolite was substantially different from that analysed in the present paper, a useful and interesting comparison can be made.

### Experimental

A single crystal of a  $2M_2$  lepidolite suitable for structural studies has been obtained from the sample, from Elba (Italy), fully described in a previous work by two of the authors. The sample also contains 1M crystals for which, as well as for the  $2M_2$  crystals, detailed optical and X-ray data were given (Franzini & Sartori, 1969).

The intensity data were collected with nickel-filtered Cu  $K\alpha$  radiation ( $\lambda = 1.5418 \text{ \AA}$ ) by means of Weissenberg photographs, using the multiple-film technique and integration process. A transparent colourless flake, plate parallel to (001), with a perfect hexagonal outline up to 1.28 mm across and 0.04 mm in thickness, was carefully chosen from among the least deformed crystals (its Weissenberg photographs showed only minor diffuse scattering and spot broadening). Five layers with  $b$  as rotation axis ( $k = 0$  to 4) were recorded; 525 independent reflexions were observed. The intensities, measured with a Nonius microdensitometer, have been corrected for Lorentz and polarization factors. The absorption correction was made by computing the transmission factor by means of a program (Alberti, 1968) based on the Monte Carlo method proposed by Alberti & Gottardi (1966) for crystals of any shape and absorption: with a linear absorption coefficient of  $115.31 \text{ cm}^{-1}$  for Cu  $K\alpha$  radiation, the transmission factors were found to vary to a great extent, ranging from 0.10 to 0.65.

### Determination and refinement of the structure

A hypothetical structure for the  $2M_2$  polymorph was proposed by two of us (Franzini & Sartori, 1969) in a previous paper concerned with the crystal data of the Elba lepidolites. According to this hypothesis, based mainly on one-dimensional data, the  $2M_2$  polymorph

would be built up by the stacking of two mica layers, labelled  $A$  and  $B$ , differing structurally by the opposite rotations of tetrahedra (deviations from hexagonal symmetry); this structural model would retain an octahedral coordination around the potassium ion, notwithstanding the ditrigonal surface symmetry, for the polymorphs derived from the stacking of mica layers rotated by an odd multiple of  $60^\circ$  (i.e.  $2M_2$ ,  $2O$  and  $6H$ ).

This hypothesis has not been substantiated by the present study based on three-dimensional data; the full structural analysis, as pointed out also by Takeda *et al.* (1971), has shown that the structural features of the lepidolites cannot be elucidated by means of a one-dimensional Fourier synthesis and that the tendency of the potassium ion to assume an octahedral coordination had been overestimated.

The trial structure was first computed on the basis of such a model of alternating  $A$  and  $B$  mica layers; the atomic coordinates were derived from those of the basic mica unit (the 1M form), predicted from the cell dimensions and composition (Donnay, Donnay & Takeda, 1964), taking into account the geometric relationships between the  $A$  and  $B$  layers pointed out by Franzini (1969). The space group chosen as the most probable for this model was  $C\bar{1}$ . Statistical distribution of aluminum and lithium over all octahedral positions was assumed.

The atomic scattering factors were taken from *International Tables for X-ray Crystallography* (1962), all the atoms being considered in their neutral state. Three cycles of full-matrix least-squares refinement of positional parameters, scale factors and isotropic temperature factors were computed, by means of the program by Busing, Martin & Levy adapted for the Crystal Structures Calculations System X-ray 63 (Stewart, 1964); the resulting unweighted  $R_1$  value (defined as  $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$ ) for the observed reflexions was 0.142. No further improvement was obtained with subsequent refinement cycles, during which the shifts of positional parameters remained smaller than the standard deviations.

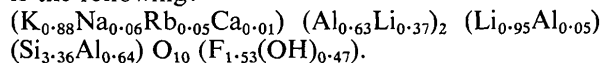
At this stage a refinement in the  $C2/c$  space group of half of the atoms refined in  $C\bar{1}$  was tried; it proved to be successful, since the  $R_1$  value rapidly dropped to 0.130. All the subsequent refinement cycles were then carried out in space group  $C2/c$  and the structural model based on alternating  $A$  and  $B$  mica layers was rejected.

A preliminary computation of the bond lengths showed that the two crystallographically independent octahedra had significantly different mean bond lengths ( $M_1-O = 1.988$ ,  $M_2-O = 2.119 \text{ \AA}$ ); this suggested a remarkable ordering in octahedral cation distribution. On the basis of the values given by Shannon & Prewitt (1969) for Li-O (2.12  $\text{\AA}$ ) and Al-O (1.91  $\text{\AA}$ ) distances for lithium and aluminum in octahedral coordination, we computed occupancies corresponding to  $Al_{0.63}Li_{0.37}$  for position 8( $f$ ) and  $Li_{0.95}Al_{0.05}$  for position 4( $c$ ). (Actually the  $M_2-O$  mean distance did not indicate any isomorphic replacement so that at first only Li was

placed in this octahedron; the negative value assumed by its temperature factor forced us to admit for this position a small amount of substitution of aluminum for lithium).

After these changes and two further cycles of refinement the  $R_1$  value was reduced to 0.117. By a careful survey of the calculated structure factors a few mistakes in estimating intensities and indexing reflexions were corrected; the final value assumed by the unweighted  $R_1$  index, for all the 525 observed reflexions, was 0.096, whereas the weighted  $R_2$  index (defined as  $[\sum w|F_o| - |F_c|]^2 / \sum w|F_o|^2$ ) reached 0.106 (in all refinement cycles unit weights were assigned to all reflexions).

No further changes of the octahedral occupancies have been made in the last cycles, since the shifts of the  $M_{\text{oct}}-O$  mean bond lengths were smaller than their standard deviations; hence the final structural formula is the following:



The observed and calculated factors are compared in Table 1; Table 2 gives the final positional and thermal parameters with their standard deviations.

### Discussion of the structure

Before discussing in detail the features of the structure, a few remarks about the model proposed by Franzini

Table 1. Observed and calculated structure factors for  $2M_2$  lepidolite

hkl	Observed	Calculated	hkl	Observed	Calculated
1 0 0	1000	1000	1 0 0	1000	1000
2 0 0	4000	4000	2 0 0	4000	4000
3 0 0	9000	9000	3 0 0	9000	9000
4 0 0	16000	16000	4 0 0	16000	16000
5 0 0	25000	25000	5 0 0	25000	25000
6 0 0	36000	36000	6 0 0	36000	36000
7 0 0	49000	49000	7 0 0	49000	49000
8 0 0	64000	64000	8 0 0	64000	64000
9 0 0	81000	81000	9 0 0	81000	81000
10 0 0	100000	100000	10 0 0	100000	100000
11 0 0	121000	121000	11 0 0	121000	121000
12 0 0	144000	144000	12 0 0	144000	144000
13 0 0	169000	169000	13 0 0	169000	169000
14 0 0	196000	196000	14 0 0	196000	196000
15 0 0	225000	225000	15 0 0	225000	225000
16 0 0	256000	256000	16 0 0	256000	256000
17 0 0	289000	289000	17 0 0	289000	289000
18 0 0	324000	324000	18 0 0	324000	324000
19 0 0	361000	361000	19 0 0	361000	361000
20 0 0	400000	400000	20 0 0	400000	400000
21 0 0	441000	441000	21 0 0	441000	441000
22 0 0	484000	484000	22 0 0	484000	484000
23 0 0	529000	529000	23 0 0	529000	529000
24 0 0	576000	576000	24 0 0	576000	576000
25 0 0	625000	625000	25 0 0	625000	625000
26 0 0	676000	676000	26 0 0	676000	676000
27 0 0	729000	729000	27 0 0	729000	729000
28 0 0	784000	784000	28 0 0	784000	784000
29 0 0	841000	841000	29 0 0	841000	841000
30 0 0	899000	899000	30 0 0	899000	899000
31 0 0	959000	959000	31 0 0	959000	959000
32 0 0	1020000	1020000	32 0 0	1020000	1020000
33 0 0	1082000	1082000	33 0 0	1082000	1082000
34 0 0	1145000	1145000	34 0 0	1145000	1145000
35 0 0	1209000	1209000	35 0 0	1209000	1209000
36 0 0	1274000	1274000	36 0 0	1274000	1274000
37 0 0	1340000	1340000	37 0 0	1340000	1340000
38 0 0	1407000	1407000	38 0 0	1407000	1407000
39 0 0	1475000	1475000	39 0 0	1475000	1475000
40 0 0	1544000	1544000	40 0 0	1544000	1544000
41 0 0	1614000	1614000	41 0 0	1614000	1614000
42 0 0	1685000	1685000	42 0 0	1685000	1685000
43 0 0	1757000	1757000	43 0 0	1757000	1757000
44 0 0	1830000	1830000	44 0 0	1830000	1830000
45 0 0	1904000	1904000	45 0 0	1904000	1904000
46 0 0	1979000	1979000	46 0 0	1979000	1979000
47 0 0	2055000	2055000	47 0 0	2055000	2055000
48 0 0	2132000	2132000	48 0 0	2132000	2132000
49 0 0	2210000	2210000	49 0 0	2210000	2210000
50 0 0	2289000	2289000	50 0 0	2289000	2289000
51 0 0	2369000	2369000	51 0 0	2369000	2369000
52 0 0	2450000	2450000	52 0 0	2450000	2450000
53 0 0	2532000	2532000	53 0 0	2532000	2532000
54 0 0	2615000	2615000	54 0 0	2615000	2615000
55 0 0	2700000	2700000	55 0 0	2700000	2700000
56 0 0	2786000	2786000	56 0 0	2786000	2786000
57 0 0	2873000	2873000	57 0 0	2873000	2873000
58 0 0	2961000	2961000	58 0 0	2961000	2961000
59 0 0	3050000	3050000	59 0 0	3050000	3050000
60 0 0	3140000	3140000	60 0 0	3140000	3140000
61 0 0	3231000	3231000	61 0 0	3231000	3231000
62 0 0	3323000	3323000	62 0 0	3323000	3323000
63 0 0	3416000	3416000	63 0 0	3416000	3416000
64 0 0	3510000	3510000	64 0 0	3510000	3510000
65 0 0	3605000	3605000	65 0 0	3605000	3605000
66 0 0	3701000	3701000	66 0 0	3701000	3701000
67 0 0	3798000	3798000	67 0 0	3798000	3798000
68 0 0	3896000	3896000	68 0 0	3896000	3896000
69 0 0	3995000	3995000	69 0 0	3995000	3995000
70 0 0	4095000	4095000	70 0 0	4095000	4095000
71 0 0	4196000	4196000	71 0 0	4196000	4196000
72 0 0	4298000	4298000	72 0 0	4298000	4298000
73 0 0	4401000	4401000	73 0 0	4401000	4401000
74 0 0	4505000	4505000	74 0 0	4505000	4505000
75 0 0	4610000	4610000	75 0 0	4610000	4610000
76 0 0	4716000	4716000	76 0 0	4716000	4716000
77 0 0	4823000	4823000	77 0 0	4823000	4823000
78 0 0	4931000	4931000	78 0 0	4931000	4931000
79 0 0	5040000	5040000	79 0 0	5040000	5040000
80 0 0	5150000	5150000	80 0 0	5150000	5150000
81 0 0	5261000	5261000	81 0 0	5261000	5261000
82 0 0	5373000	5373000	82 0 0	5373000	5373000
83 0 0	5486000	5486000	83 0 0	5486000	5486000
84 0 0	5600000	5600000	84 0 0	5600000	5600000
85 0 0	5715000	5715000	85 0 0	5715000	5715000
86 0 0	5831000	5831000	86 0 0	5831000	5831000
87 0 0	5948000	5948000	87 0 0	5948000	5948000
88 0 0	6066000	6066000	88 0 0	6066000	6066000
89 0 0	6185000	6185000	89 0 0	6185000	6185000
90 0 0	6305000	6305000	90 0 0	6305000	6305000
91 0 0	6426000	6426000	91 0 0	6426000	6426000
92 0 0	6548000	6548000	92 0 0	6548000	6548000
93 0 0	6671000	6671000	93 0 0	6671000	6671000
94 0 0	6795000	6795000	94 0 0	6795000	6795000
95 0 0	6920000	6920000	95 0 0	6920000	6920000
96 0 0	7046000	7046000	96 0 0	7046000	7046000
97 0 0	7173000	7173000	97 0 0	7173000	7173000
98 0 0	7301000	7301000	98 0 0	7301000	7301000
99 0 0	7430000	7430000	99 0 0	7430000	7430000
100 0 0	7560000	7560000	100 0 0	7560000	7560000
101 0 0	7691000	7691000	101 0 0	7691000	7691000
102 0 0	7823000	7823000	102 0 0	7823000	7823000
103 0 0	7956000	7956000	103 0 0	7956000	7956000
104 0 0	8090000	8090000	104 0 0	8090000	8090000
105 0 0	8225000	8225000	105 0 0	8225000	8225000
106 0 0	8361000	8361000	106 0 0	8361000	8361000
107 0 0	8498000	8498000	107 0 0	8498000	8498000
108 0 0	8636000	8636000	108 0 0	8636000	8636000
109 0 0	8775000	8775000	109 0 0	8775000	8775000
110 0 0	8915000	8915000	110 0 0	8915000	8915000
111 0 0	9056000	9056000	111 0 0	9056000	9056000
112 0 0	9198000	9198000	112 0 0	9198000	9198000
113 0 0	9341000	9341000	113 0 0	9341000	9341000
114 0 0	9485000	9485000	114 0 0	9485000	9485000
115 0 0	9630000	9630000	115 0 0	9630000	9630000
116 0 0	9776000	9776000	116 0 0	9776000	9776000
117 0 0	9923000	9923000	117 0 0	9923000	9923000
118 0 0	10071000	10071000	118 0 0	10071000	10071000
119 0 0	10220000	10220000	119 0 0	10220000	10220000
120 0 0	10370000	10370000	120 0 0	10370000	10370000
121 0 0	10520000	10520000	121 0 0	10520000	10520000
122 0 0	10670000	10670000	122 0 0	10670000	10670000
123 0 0	10820000	10820000	123 0 0	10820000	10820000
124 0 0	10970000	10970000	124 0 0	10970000	10970000
125 0 0	11120000	11120000	125 0 0	11120000	11120000
126 0 0	11270000	11270000	126 0 0	11270000	11270000
127 0 0	11420000	11420000	127 0 0	11420000	11420000
128 0 0	11570000	11570000	128 0 0	11570000	11570000
129 0 0	11720000	11720000	129 0 0	11720000	11720000
130 0 0	11870000	11870000	130 0 0	11870000	11870000
131 0 0	12020000	12020000	131 0 0	12020000	12020000
132 0 0	12170000	12170000	132 0 0	12170000	12170000
133 0 0	12320000	12320000	133 0 0	12320000	12320000
134 0 0	12470000	12470000	134 0 0	12470000	12470000
135 0 0	12620000	12620000	135 0 0	12620000	12620000
136 0 0	12770000	12770000	136 0 0	12770000	12770000
137 0 0	12920000	12920000	137 0 0	12920000	12920000
138 0 0	13070000	13070000	138 0 0	13070000	13070000
139 0 0	13220000	13220000	139 0 0	13220000	13220000
140 0 0	13370000	13370000	140 0 0	13370000	13370000
141 0 0	13520000	13520000	141 0 0	13520000	13520000
142 0 0	13670000	13670000	142 0 0	13670000	13670000
143 0 0	13820000	13820000	143 0 0	13820000	13820000
144 0 0	13970000	13970000	144 0 0	13970000	13970000
145 0 0	14120000	14120000	145 0 0	14120000	14120000
146 0 0	14270000	14270000	146 0 0	14270000	14270000
147 0 0	14420000	14420000	147 0 0	14420000	14420000
148 0 0	14570000	14570000	148 0 0	14570000	14570000
149 0 0	14720000	14720000	149 0 0	14720000	14720000
150 0 0	14870000	14870000	150 0 0	14870000	14870000
151 0 0	15020000	15020000	151 0 0	15020000	15020000
152 0 0	15170000	15170000	152 0 0	15170000	15170000
153 0 0	15320000	15320000</			

From Fig. 1 and Tables 3 and 4, in which the more interesting interatomic distances and angles are reported, the fundamental features of the structure of this polymorph can be deduced.

In the tetrahedral sheet there is no evidence of any cation ordering, since the two crystallographically inde-

Table 3 Interatomic distances in 2M<sub>2</sub> lepidolite

T <sub>1</sub> tetrahedron		T <sub>2</sub> tetrahedron	
T <sub>1</sub> -O <sub>4</sub>	1.634 (13) Å	T <sub>2</sub> -O <sub>4</sub>	1.640 (13) Å
T <sub>1</sub> -O <sub>6</sub>	1.642 (11)	T <sub>2</sub> -O <sub>5</sub>	1.628 (14)
T <sub>1</sub> -O <sub>5</sub> †	1.640 (15)	T <sub>2</sub> -O <sub>6</sub> ‡	1.629 (12)
T <sub>1</sub> -O <sub>1</sub> *	1.604 (10)	T <sub>2</sub> -O <sub>2</sub> *	1.612 (10)
Mean	1.630 (6)	Mean	1.627 (6)
O <sub>4</sub> -O <sub>6</sub>	2.618 (17)	O <sub>4</sub> -O <sub>5</sub>	2.614 (21)
O <sub>6</sub> -O <sub>5</sub> <sup>†</sup>	2.642 (18)	O <sub>5</sub> -O <sub>6</sub> <sup>‡</sup>	2.616 (17)
O <sub>5</sub> -O <sub>4</sub>	2.640 (21)	O <sub>6</sub> <sup>†</sup> -O <sub>4</sub>	2.635 (17)
O <sub>6</sub> -O <sub>1</sub> *	2.678 (16)	O <sub>4</sub> -O <sub>2</sub> *	2.689 (15)
O <sub>5</sub> -O <sub>1</sub> *	2.694 (16)	O <sub>5</sub> -O <sub>2</sub> *	2.698 (16)
O <sub>4</sub> -O <sub>1</sub> *	2.688 (16)	O <sub>6</sub> <sup>†</sup> -O <sub>2</sub> *	2.684 (15)
Mean	2.660 (7)	Mean	2.656 (7)

M <sub>1</sub> octahedron		M <sub>2</sub> octahedron	
M <sub>1</sub> <sup>iv</sup> -O <sub>1</sub>	1.981 (13) Å	M <sub>2</sub> -O <sub>2</sub>	2.141 (12) Å
M <sub>1</sub> <sup>iv</sup> -O <sub>2</sub> <sup>iii</sup>	1.977 (13)	M <sub>2</sub> -O <sub>1</sub> <sup>†</sup>	2.139 (14)
M <sub>1</sub> <sup>iv</sup> -O <sub>1</sub> <sup>v</sup>	1.975 (13)	M <sub>2</sub> -O <sub>3</sub>	2.090 (11)
M <sub>1</sub> <sup>iv</sup> -O <sub>2</sub> <sup>v</sup>	1.954 (15)	Mean	2.123 (7)
M <sub>1</sub> <sup>iv</sup> -O <sub>3</sub> †	2.005 (16)		
M <sub>1</sub> <sup>iv</sup> -O <sub>3</sub> <sup>v</sup>	1.990 (13)		
Mean	1.980 (6)		

Unshared		Unshared	
O <sub>1</sub> -O <sub>2</sub> <sup>iii</sup>	2.929 (15)	O <sub>3</sub> -O <sub>2</sub>	3.214 (15)
O <sub>1</sub> -O <sub>3</sub>	2.934 (20)	O <sub>1</sub> <sup>†</sup> -O <sub>3</sub>	3.211 (20)
O <sub>3</sub> -O <sub>2</sub> <sup>iii</sup>	2.926 (20)	O <sub>2</sub> -O <sub>1</sub> <sup>†</sup>	3.205 (20)
O <sub>1</sub> <sup>v</sup> -O <sub>2</sub> <sup>v</sup>	2.916 (20)	Mean	3.210 (11)
O <sub>1</sub> <sup>v</sup> -O <sub>3</sub> <sup>v</sup>	2.908 (15)		
O <sub>3</sub> -O <sub>2</sub> <sup>v</sup>	2.914 (20)		
Mean	2.921 (7)		

Shared	
O <sub>2</sub> <sup>vii</sup> -O <sub>3</sub>	2.753 (17) Å
O <sub>1</sub> <sup>v</sup> -O <sub>3</sub>	2.753 (16)
O <sub>1</sub> <sup>v</sup> -O <sub>2</sub>	2.837 (16)
Mean	2.781 (9)
O <sub>3</sub> -O <sub>2</sub> <sup>v</sup>	2.508 (15)

Inner		Outer	
K <sup>viii</sup> -O <sub>6</sub>	2.965 (16) Å	K <sup>viii</sup> -O <sub>4</sub>	3.214 (15) Å
K <sup>viii</sup> -O <sub>5</sub>	2.981 (15)	K <sup>viii</sup> -O <sub>6</sub> <sup>†</sup>	3.250 (17)
K <sup>viii</sup> -O <sub>4</sub> <sup>viii</sup>	2.982 (14)	K <sup>viii</sup> -O <sub>5</sub> <sup>iii</sup>	3.320 (15)
Mean	2.976 (9)	Mean	3.261 (9)
K <sup>viii</sup> -O <sub>3</sub>		3.997 (10) Å	

Shortest interlayer distances	
O <sub>6</sub> -O <sub>6</sub> <sup>ix</sup>	3.368 (14) Å (×2)
O <sub>5</sub> <sup>iii</sup> -O <sub>5</sub> <sup>ix</sup>	3.430 (15) (×4)
Mean	3.409 (8)

\* Apical oxygens

† O<sub>3</sub> Hydroxyl and fluorine

‡ The atoms of the different asymmetric units are related to the symmetry-equivalent atoms of the fundamental unit as follows:

i	x, y-1, z	v	1-x, 1-y, -z
ii	x-½, y+½, z	vi	x, 1+y, z
iii	x+½, y-½, z	vii	½-x, ½-y, -z
iv	½-x, ½-y, -z	viii	½+x, ½+y, z
		ix	1-x, y, ½-z

Table 4. Interatomic angles in 2M<sub>2</sub> lepidolite  
Central atom is the vertex

T <sub>1</sub> tetrahedron		T <sub>2</sub> tetrahedron	
O <sub>5</sub> <sup>†</sup> -T <sub>1</sub> -O <sub>6</sub> †	107.20 (64)°	O <sub>6</sub> <sup>†</sup> -T <sub>2</sub> -O <sub>5</sub>	106.89 (67)°
O <sub>5</sub> <sup>†</sup> -T <sub>1</sub> -O <sub>4</sub>	107.50 (68)	O <sub>6</sub> <sup>†</sup> -T <sub>2</sub> -O <sub>4</sub>	107.40 (65)
O <sub>6</sub> <sup>†</sup> -T <sub>1</sub> -O <sub>4</sub>	106.09 (63)	O <sub>5</sub> <sup>†</sup> -T <sub>2</sub> -O <sub>4</sub>	106.19 (65)
O <sub>1</sub> <sup>*</sup> -T <sub>1</sub> -O <sub>6</sub>	111.18 (63)	O <sub>2</sub> <sup>*</sup> -T <sub>2</sub> -O <sub>4</sub>	111.53 (66)
O <sub>1</sub> <sup>*</sup> -T <sub>1</sub> -O <sub>4</sub>	112.23 (64)	O <sub>2</sub> <sup>*</sup> -T <sub>2</sub> -O <sub>5</sub>	112.71 (66)
O <sub>1</sub> <sup>*</sup> -T <sub>1</sub> -O <sub>5</sub>	112.29 (66)	O <sub>2</sub> <sup>*</sup> -T <sub>2</sub> -O <sub>6</sub> <sup>†</sup>	111.77 (58)
Mean	109.42 (26)	Mean	109.42 (26)

Angles between basal oxygens:		Amount of rotation from 120°
O <sub>6</sub> -O <sub>4</sub> -O <sub>5</sub>	107.28 (57)°	12.72°
O <sub>4</sub> -O <sub>5</sub> -O <sub>6</sub> <sup>†</sup>	132.80 (60)	12.80
O <sub>5</sub> -O <sub>6</sub> <sup>†</sup> -O <sub>4</sub> <sup>iii</sup>	106.85 (66)	13.15
O <sub>5</sub> <sup>†</sup> -O <sub>4</sub> <sup>iii</sup> -O <sub>6</sub> <sup>†</sup>	132.55 (59)	12.55
O <sub>4</sub> <sup>iii</sup> -O <sub>5</sub> <sup>†</sup> -O <sub>6</sub>	106.98 (58)	13.02
O <sub>5</sub> <sup>†</sup> -O <sub>6</sub> -O <sub>4</sub>	133.23 (70)	13.23
Mean		12.91 (25)
Tetrahedral rotation angle α = 6.45 (13)°.		

\* Apical oxygens.

† The atoms of the different asymmetric units are related to the symmetry equivalent atoms of the fundamental unit as follows:

i	x, y-1, z	v	1-x, 1-y, -z
ii	x-½, y+½, z	vi	x, 1+y, z
iii	x+½, y-½, z	vii	½-x, ½-y, -z
iv	½-x, ½-y, -z	viii	½+x, ½+y, z
		ix	1-x, y, ½-z

pendent tetrahedra are identical within the precision of the determination. The tetrahedra are rather small in comparison with those of all the other micas, both

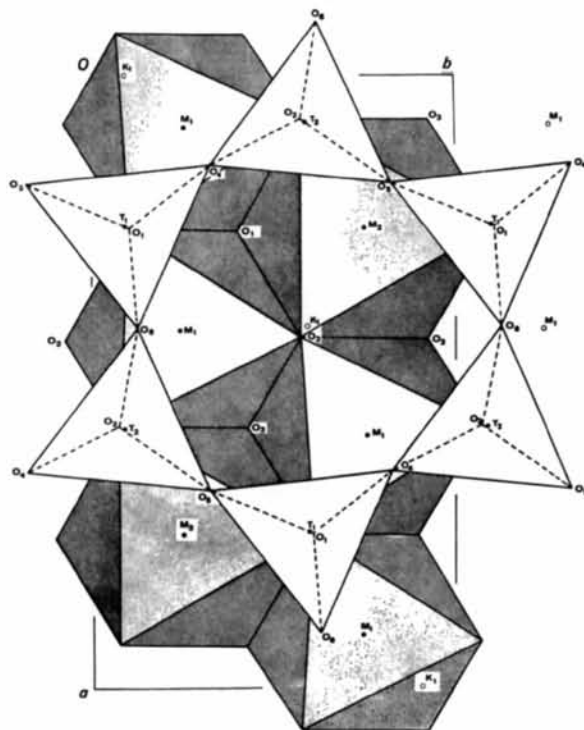


Fig. 1. View along c\* of the 2M<sub>2</sub> lepidolite structure showing one (Si<sub>2</sub>O<sub>5</sub>)<sup>2-</sup> ring and part of an octahedral sheet.

di octahedral and trioctahedral, because of the little replacement of Al for Si. Moreover, as we have already emphasized, these coordination polyhedra are remarkably distorted into a form near to the trigonal pyramid.

Noteworthy is the shortening of the Si–O<sub>apex</sub> distance (mean length 1.608 Å), compared with the Si–O<sub>base</sub> distance (mean length 1.635 Å), a peculiarity which is also encountered in the 2M<sub>2</sub> lepidolite structure determined by Takeda *et al.* (1971), and it is even more evident in the polyolithionite structure. This behaviour can be explained on the basis of the *d*–*p*  $\pi$  bond theory suggested by Cruickshank (1961) as well as on the basis of the Pauling–Zachariasen method of the balancing of valences. If appreciable  $\pi$  bonding is responsible for the short Si–O<sub>apex</sub> bond, the bond angles subtended at silicon by the tetrahedral edges and involving O<sub>apex</sub> should be affected. They are indeed, with angles 111 to 113°, whereas the angles subtended by the edges made of bridging oxygen atoms range from 106 to 107.5° (Table 4).

In the octahedral sheet we find a considerable ordering of the cations; in fact the lithium ions almost completely fill the 4(*c*) position before entering the 8(*f*) position. Consequently the octahedral sheet appears to be built up from a large octahedron and two smaller ones, assuming a configuration rather similar to that of the corresponding sheet of the dioctahedral micas, with the difference that whereas the Li-rich octahedron is rather close in size to an empty octahedron, the other two octahedra, which still contain an appreciable amount of lithium (Al<sub>0.63</sub>Li<sub>0.37</sub>), are larger than the filled octahedra of the dioctahedral micas. As a result the lateral dimensions of the octahedral sheet are rather large, also owing to the flattening of such coordination polyhedra (mainly of the Li-rich octahedron,  $\psi = 60^\circ 47'$ ).

It is then evident that the misfit between such an octahedral sheet and a tetrahedral sheet characterized by small elongated tetrahedra is quite small, and consequently the tetrahedral rotation needed for the adjustment of one sheet to the other is also small. The 'tetrahedral tilt' (calculated according to the formula:  $\Delta z = [(ZO_6 + ZO_4)/2 - ZO_5] \cdot c \sin \beta$ ) is also small; the value  $\Delta z = 0.086$  attests that the O<sub>basal</sub> sheet is, on the whole, little corrugated.

As in other micas the K–O distances are of two kinds, due to the reduced, but always manifest, ditrigonal configuration of the oxygen rings on either side of the potassium ion. The observed values (Table 3) give averages of 2.976 and 3.261 Å for the inner and the outer K–O distances, respectively. Since the K–O distances depend on the size of the tetrahedra and on their rotation angle (Takeda & Burnham, 1966), the small size of the tetrahedra and the low value of  $\alpha$  account for the small value of the longer of the two K–O distances and explain why the shorter K–O distance is as large as it is. Since in 2M<sub>2</sub> polymorph the two mica layers are rotated by 60° on each other, the coordination polyhedron around K<sup>+</sup> is then a trigonal

prism if we take into account only the 6 inner oxygens, a ditrigonal prism if we consider also the 6 outer oxygens.

A detailed comparison of the parameters of our mica with those of the other two lepidolites whose structures are reported in the literature now seems appropriate. Most interesting, in our opinion, is a comparison with Takeda's 2M<sub>2</sub> lepidolite structure. Tables 5 and 6 show in fact a close analogy between the structure of the Elba lepidolite and that of the Rozna lepidolite. However, these crystals are substantially different in composition; if we express their composition in terms of the three end members, polyolithionite (Pl) [KLi<sub>2</sub>AlSi<sub>4</sub>O<sub>10</sub>F<sub>2</sub>], trilithionite (Tl) [KLi<sub>3/2</sub>Al<sub>3/2</sub>Si<sub>3</sub>AlO<sub>10</sub>(F,OH)<sub>2</sub>] and muscovite (Ms) [KAl<sub>2</sub>Si<sub>3</sub>AlO<sub>10</sub>(OH)<sub>2</sub>], the composition of our lepidolite is close to Pl<sub>50</sub>Tl<sub>50</sub>, whereas that of the Takeda's mica is Ms<sub>40</sub>Pl<sub>30</sub>Tl<sub>30</sub>. These compositions, plotted in the ternary diagram of Munoz (1968), turn out to be of great interest, since Takeda's mica is very close to the region of the mixed-layer Li-micas, while our mica lies on the Pl–Tl join, in the field of the 1M and 2M<sub>2</sub> lepidolites.

Table 5. Cell dimensions and structural formulae of 2M<sub>2</sub> lepidolites

	Lepidolite, 2M <sub>2</sub> *	Lepidolite, 2M <sub>2</sub> †
<i>a</i>	9.04 (2) Å	9.032 (2) Å
<i>b</i>	5.22 (2)	5.200 (3)
<i>c</i>	20.210 (1)	20.15 (4)
$\beta$	99° 35' (20')	99° 46' (10')
Space group	C2/c	C2/c

\* (K<sub>0.88</sub>Na<sub>0.06</sub>Rb<sub>0.05</sub>Ca<sub>0.01</sub>)(Al<sub>0.63</sub>Li<sub>0.37</sub>)<sub>2</sub>(Li<sub>0.95</sub>Al<sub>0.05</sub>)(Si<sub>3.36</sub>Al<sub>0.64</sub>)O<sub>10</sub>(F<sub>1.53</sub>(OH)<sub>0.47</sub>): Elba, Italy. Present work.

† (K<sub>0.87</sub>Na<sub>0.12</sub>Rb<sub>0.06</sub>Ca<sub>0.02</sub>)(Al<sub>0.65</sub>Li<sub>0.35</sub>)<sub>2</sub>(Li<sub>0.35</sub>Al<sub>0.10</sub>(Fe<sup>2+</sup>, Mn, Mg)<sub>0.15</sub>□<sub>0.40</sub>)(Si<sub>3.39</sub>Al<sub>0.61</sub>)O<sub>10</sub>(F<sub>1.2</sub>(OH)<sub>0.8</sub>): Rozna, Moravia, Czechoslovakia. Takeda *et al.* (1971).

Table 6. Comparison of some interatomic distances in 2M<sub>2</sub> lepidolites

	2M <sub>2</sub> lepidolite from Elba (Present work)	2M <sub>2</sub> lepidolite from Rozna, Moravia (Takeda <i>et al.</i> , 1971)
	Tetrahedral sheet	
Mean T <sub>1</sub> –O	1.630 (6) Å	1.622 (4) Å
Mean T <sub>2</sub> –O	1.627 (6)	1.633 (4)
	Octahedral sheet	
Mean M <sub>1</sub> –O	1.980 (6) Å	1.967 (3) Å
Mean M <sub>2</sub> –O	2.123 (7)	2.144 (4)
	Interlayer region	
Mean K–O (inner)	2.976 (9) Å	2.980 (5) Å
Mean K–O (outer)	3.261 (9)	3.220 (5)

These compositions are calculated on the basis of chemical analyses of samples which contain both the 1M and 2M<sub>2</sub> polymorphs. The doubt arises that such compositions are not necessarily related to the 2M<sub>2</sub> polymorphs structurally analysed and that the close resemblance between the structural parameters re-

ported by Takeda *et al.* and ours is to be ascribed to the true chemical compositions of the crystals being much closer to each other than the bulk compositions of the samples. Though lacking chemical data on a single crystal, we believe that the satisfactory refinements of the structures rule out this possibility; the reason for the close resemblance of the two structures is then to be found in other features. We think such a reason most likely lies in the following two peculiarities.

The two  $2M_2$  lepidolites have tetrahedral sheets nearly equal in composition ( $Si_{3.39}Al_{0.61}$  = Rozna lepidolite;  $Si_{3.36}Al_{0.64}$  = Elba lepidolite) and so are very close in dimensions.

They have octahedral sheets which, though clearly different in composition, are both made up of a large octahedron and two smaller ones, whose sizes are very close in the two lepidolites. Such a close geometric resemblance of the octahedral sheets of the two crystals is achieved through an almost equal composition of the small octahedra ( $Al_{0.65}Li_{0.35}$  = Rozna lepidolite;  $Al_{0.63}Li_{0.37}$  = Elba lepidolite) and through the concentration in the larger octahedron of the Rozna lepidolite of all the vacancies, while the corresponding site of the Elba lepidolite accommodates almost only lithium.

The  $1M$  polyolithionite studied by Takeda & Burnham (1969), in spite of its different polymorphic type, has tetrahedral and octahedral sheets which show features rather close to those of the  $2M_2$  lepidolites. They are even more evident because of the presence of only Si in the first sheet and of the higher content of Li of the second; as a consequence the misfit between them is even less and the tetrahedral rotation angle  $\alpha$  lowest ( $\alpha = 3^\circ$ ).

To sum up, we can say that the most interesting feature which emerges from the comparison between all the structurally analysed lepidolites is the constant presence of remarkable octahedral ordering. The Li-Al substitution is characterized by concentrating in an octahedral site almost exclusively Li (0.90 to 0.95 in polyolithionite and in the Elba lepidolite) or Li plus vacancies (0.75 in the Rozna lepidolite), and by filling the other two sites with Al plus the remaining Li. It is in this distorted octahedral sheet and in its consequently easy adaptability to the tetrahedral sheet without a great deviation from hexagonality of the latter, that we think lies the main difference between the structure of the polyolithionite-rich lepidolites and that of the lepidolites

close to trilithionite, characterized by having the  $2M_1$  as the more stable polymorph (Munoz, 1968). The discontinuity observed by Munoz (1968) between them would then be a discontinuity which separates a structure with a remarkable octahedral ordering from one with a disordered octahedral sheet. Rieder (1968), in his paper on lithium-iron micas, also suggested that the octahedral sheet of the trilithionites must be disordered.

However the final answer to the problem of the relations between the structure of trilithionites and that of the other lepidolites and between the former and that of the  $2M_1$  lithian muscovites can be given only by the crystal structure analysis of a  $2M_1$  trilithionite.

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